# Evolution of magnetism in $U(Ni_{1-x}Pd_x)_2Si_2$ single crystals

F. Honda<sup>1</sup>, A.V. Andreev<sup>2</sup>, V. Sechovský<sup>1,a</sup>, Y. Homma<sup>3</sup>, and Y. Shiokawa<sup>3</sup>

<sup>1</sup> Department of Electronic Structures, Charles University, Ke Karlovu 5, 121 16 Prague 2, The Czech Republic

<sup>2</sup> Institute of Physics ASCR, Na Slovance 2, 182 21 Prague 8, The Czech Republic

<sup>3</sup> Institute for Materials Research, Tohoku University, Katahira 2-1-1, Aoba-ku, Sendai 980-8577, Japan

Received 28 March 2002 / Received in final form 8 August 2002 Published online 19 December 2002 – © EDP Sciences, Società Italiana di Fisica, Springer-Verlag 2002

**Abstract.** Single crystals of  $U(Ni_{1-x}Pd_x)_2Si_2$  with x = 0.05, 0.09 and 0.135 have been grown. Magnetization and electrical resistivity measurements were performed in a wide range of temperatures and magnetic fields in order to study stability of magnetic phases in the solid solutions between  $UNi_2Si_2$  and  $UPd_2Si_2$  with a special emphasis on the type of ground state. In  $UPd_2Si_2$  the simple AFI-type antiferromagnetic structure of U moments is observed at low temperatures.  $UNi_2Si_2$  adopts the uncompensated AF structure (UAF) with the ++- stacking of U moments along the *c*-axis and consequently this compound exhibits a spontaneous magnetization corresponding to 1/3 of the U moment. The substitution of Pd for Ni leads to a rapid decay of the spontaneous magnetization. The evolution of magnetization and electrical resistivity behavior with Pd doping is tentatively attributed to the coexistence of the AF-I and UAF phases in the ground state of  $U(Ni_{0.91}Pd_{0.09})_2Si_2$  and  $U(Ni_{0.865}Pd_{0.135})_2Si_2$ . In this scenario, the volume fraction of the AF-I phase rapidly grows with Pd doping on account of the UAF. At lowest temperatures an irreversible transition to the UAF phase is observed when a sufficiently high magnetic field is applied along the *c*-axis.

**PACS.** 75.50 Ee Antiferromagnetic materials – 75.50 Gg Ferrimagnetics – 75.30 Gw Magnetic anisotropy – 75.30 Kz Magnetic phase transitions

# 1 Introduction

The  $UT_2Si_2$  (T: transition metals) compounds crystallize in the body-centered ThCr<sub>2</sub>Si<sub>2</sub>-type of tetragonal structure characterized by the space group I4/mmm [1]. The crystal structure consists of the U basal planes, which are separated by the T-Si slabs. The  $UT_2Si_2$  compounds with various transition metals exhibit variety of interesting physical properties such as the heavy-fermion state, superconductivity, 5f-electron magnetism or coexistence of the two cooperative phenomena [2]. Closer inspection of experimental data collected for the  $UT_2Si_2$  family reveals that the compounds with the transition metals from the same column of the Periodic table frequently exhibit common features. Both, the UPd<sub>2</sub>Si<sub>2</sub> and UNi<sub>2</sub>Si<sub>2</sub> compounds order magnetically at relatively high temperatures  $(T_{\rm N} = 133 \text{ and } 124 \text{ K}, \text{ respectively})$  [3,4]. In a limited temperature range (down to  $T_2 = 108$  and 103 K, respectively) both compounds exhibit an incommensurate magnetic structure of U magnetic moments, which is frequently quoted as the longitudinal spin density wave

(LSDW). At  $T_{\rm 2}$  the both compounds undergo the first-order magnetic phase transition to the body-centered an-

tiferromagnetic AF-I structure characterized by the simple +-+-+- coupling of the U magnetic moments along the c-axis. In UPd<sub>2</sub>Si<sub>2</sub> this magnetic structure persists down to lowest temperatures and is characteristic for the ground state. In UNi<sub>2</sub>Si<sub>2</sub> the AF-I phase is stable only in a limited temperature range and at  $T_1 = 53$  K the AF-I phase transforms to the ground-state phase, which is characterized by a spontaneous magnetization.

Magnetism in UPd<sub>2</sub>Si<sub>2</sub> and UNi<sub>2</sub>Si<sub>2</sub>, as well as in the other  $UT_2Si_2$  counterparts is characterized by the strong uniaxial anisotropy with the easy-magnetization direction along the *c*-axis. On the microscopic level this anisotropy is associated with the ordered U magnetic moments being locked always along c in this ThCr<sub>2</sub>Si<sub>2</sub>-type structure. All the magnetic structures observed in the two compounds have in common basal-plane sheets of ferromagnetically coupled U moments whereas they differ in the propagation of the U moment along the c-axis; *i.e.* all the propagation vectors have the general form q = (0,0,q). The LSDW phase in  $UNi_2Si_2$  is described by the temperature dependent value of  $q \sim 0.74$  [4,5] whereas the value of 0.732 was determined for UPd<sub>2</sub>Si<sub>2</sub> [3]. The AF-I phase is characterized by q = 1 whereas the ground-state phase in  $UNi_2Si_2$ , which exhibits the spontaneous moment equal

<sup>&</sup>lt;sup>a</sup> e-mail: sech@mag.mff.cuni.cz

to 1/3 of the U moment as a result of the ++- stacking, is quoted as the uncompensated antiferromagnetic phase (UAF) with q = 2/3 [4,5]. A detailed analysis of magnetization data revealed that besides the main AF-I phase the low-temperature state in UPd<sub>2</sub>Si<sub>2</sub> contains about 1% of the UAF phase. The results of magnetization measurements at 1.3 K [6] indicate a metamagnetic transition into the UAF phase in a field of ~ 15 T applied along the *c*-axis.

The magnetic phase transitions of  $UPd_2Si_2$  and  $UNi_2Si_2$  are accompanied by clear anomalies in the electrical resistivity [7,8], which makes the resistivity measurements a useful tool to study magnetism in these materials.

The stability of the AF-I and UAF phases with respect to the concentration of Ni and Pd is a subject of the present research. The pilot study performed on polycrystals of the  $U(Ni_{1-x}Pd_x)_2Si_2$  solid solutions led to the conclusion that in compounds for  $x \ge 0.25$  the ground-state phase is characterized by the entire AF-I phase [9]. In the present work, we have prepared several  $U(Ni_{1-x}Pd_x)_2Si_2$  single crystals with various Pd content  $x \le 0.135$  and performed magnetization and electrical resistivity measurements. The main motivation was to investigate the evolution of the ground state for the Ni/Pd concentration region, for which the transformation between the AF-I and UAF phase is expected.

#### 2 Experimental

The U(Ni<sub>1-x</sub>Pd<sub>x</sub>)<sub>2</sub>Si<sub>2</sub> single crystals used for the present study were grown in a tetra-arc furnace by a modified Czochralski method from initial mixtures of components corresponding to x = 0.05, 0.10 and 0.15. Phase purity was confirmed by X-ray powder diffraction and by a microprobe analysis. The microprobe analysis confirmed the exact 1:2:2 stoichiometry for all the compounds studied (within a 1% error). As regards to the Pd-Ni sublattice, the crystals were found to correspond to average x = 0.05, 0.09 and 0.135, respectively. The composition varies over the crystals within  $\Delta x = 0.01$  for x = 0.05 and 0.09 and  $\Delta x = 0.015$  for x = 0.135. The lattice parameters determined as a = 396.5 pm, c = 953.0 pm for x = 0.05; a = 397.6 pm, c = 956.7 pm for x = 0.09; a = 398.7 pm,c = 959.4 pm for x = 0.135 are in satisfactory agreement with the expected linear increase of lattice parameters between  $UNi_2Si_2$  and  $UPd_2Si_2$  [2]. The crystal quality was checked and crystal orientation was facilitated by means of X-ray Laue method.

The magnetization curves were measured at various temperatures (4.2–130 K) using a vibrating sample magnetometer with a superconducting magnet in fields up to 12 T with continuous field sweeps at a rate  $dB/dt = \pm 0.75$  T/min.

A SQUID magnetometer with magnetic fields up to 4 T was employed for measurements of the temperature dependence of the magnetization (magnetic susceptibility) in the temperature interval 5-300 K. The electrical resistivity for current along the *c*-axis was measured in the temperature range 3-300 K by a conventional four-probe AC method.



Fig. 1. The magnetization isotherms along the *c*-axis of the  $U(Ni_{0.95}Pd_{0.05})_2Si_2$  crystal at 4.2 K and elevated temperatures. The magnetization curve along the *a*-axis at 4.2 K is also shown.

### 3 Results and discussion

# 3.1 U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub>

The magnetization curves measured on U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub> in magnetic fields applied along the *c*-axis are shown in Figure 1 for several representative temperatures. At temperatures up to 40 K one observes a nearly rectangular hysteresis loop with an invariable remanent magnetic moment of 0.55  $\mu_{\rm B}$ . The value of coercive force is gradually decreasing with increasing temperature (0.8 T at 4.2 K, 0.1 T at 40 K). This magnetization behavior is identical with that of UNi<sub>2</sub>Si<sub>2</sub> [2,10], *i.e.* the UAF phase is characteristic for the ground state of both, the UNi<sub>2</sub>Si<sub>2</sub> and U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub>.

The magnetization curves measured at 60, 80 and 100 K are compatible with the existence of the AF-I phase at these temperatures. Only these typical curves are shown for clear illustration although measurements at several other temperatures between 60 and 100 K were performed, as well. These M(B) curves show no spontaneous magnetic moment but exhibit a sharp increase at a critical field  $B_c$  at which the AF-I structure transforms to the UAF phase. This metamagnetic transition exhibits a pronounced hysteresis. The value of  $B_c$  reaches maximum at temperatures around 80 K while the hysteresis becomes strongly reduced.

The magnetization curves measured at 110 and 120 K are considerably different from those observed at lower temperatures. This is because the zero-field state is no more characterized by the AF-I but it is the LSDW phase. For this temperature range, the low-field slope of the M(B) curve is rather high and the metamagnetic transition becomes broadened.

In Figure 1 also the magnetization curve is shown as measured on  $U(Ni_{0.95}Pd_{0.05})_2Si_2$  at 4.2 K in a magnetic field applied along the *a*-axis. This *M* vs. *B* dependence is



Fig. 2. a). Temperature dependence of magnetization along the *c*-axis of the x = 0.05 sample for ZFC and FC processes measured in a 0.01 T magnetic field. Expanded M(T) curves around  $T_2$  and  $T_N$  are in inset. b) Temperature dependence of magnetization along the *c*-axis of the sample with x = 0.09and 0.135.

not far from being linear (except the initial part which is sensitive to a possible small misorientation of the crystal) with a very low susceptibility and is nearly identical with the basal-plane magnetization data reported for the parent compounds UNi<sub>2</sub>Si<sub>2</sub> [2] and UPd<sub>2</sub>Si<sub>2</sub> [6,10]. Moreover, we obtained almost the same *a*-axis magnetization data also for the other two crystals  $U(Ni_{0.91}Pd_{0.09})_2Si_2$  and  $U(Ni_{0.865}Pd_{0.135})_2Si_2$ . This confirms that the huge uniaxial magnetocrystalline anisotropy of parent compounds persists also in the  $U(Ni_{1-x}Pd_x)_2Si_2$  solid solutions. Since the essential features of magnetism are concentrated to the magnetization along the *c*-axis, the following magnetization data will be presented and discussed only for magnetic fields applied along this direction.

Figure 2a displays the variation of the magnetization of U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub> in the field B = 0.01 T with increasing temperature. Note that the ZFC (FC) curve for all three compounds was measured on heating after cooling the sample in zero magnetic field (in the magnetic field of measurement). The ZFC data exhibit negligible magnetization ( $M < 10^{-3} \mu_{\rm B}/{\rm f.u.}$ ) except for the temperature interval between 25 and 50 K, where it shows a peak reaching  $M \simeq 0.05 \mu_{\rm B}/{\rm f.u.}$  at 38 K. The peak may be associated with the UAF  $\rightarrow$  AF-I magnetic phase transition. Another sharp anomaly although much weaker ( $M \simeq 4 \times 10^{-4} \mu_{\rm B}/{\rm f.u.}$ ) peaks at 105.5 K (see inset of Fig. 2a) and marks the AF-I  $\rightarrow$  LSDW magnetic phase transitions. No anomaly, however, shows up in the ZFC curve at  $T_{\rm N}$ .

The low temperature FC-magnetization data up to approximately 25 K show a high constant value of 0.41  $\mu_{\rm B}/{\rm f.u.}$  At temperatures increasing above 25 K the magnetization precipitously decreases and joins the ZFC



Fig. 3. Temperature dependence of electrical resistivity along the *c*-axis of the  $U(Ni_{1-x}Pd_x)_2Si_2$  crystals with x = 0.0 - 0.135 for heating and cooling.

curve at 38 K. The FC and ZFC curves for temperatures above 38 K are identical except for the temperature interval between  $T_2$  and  $T_N$ , where the FC curve shows somewhat higher values, which is connected with a tiny step at  $T_N$  as seen in the inset of Figure 2a.

As can be seen in Figure 3, the three magnetic phase transitions revealed by the magnetization measurements on U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub> are accompanied by corresponding anomalies on the temperature dependence of electrical resistivity. The transitions at  $T_1$  ( $T_2$ ) are associated with a rapid decrease (increase) of resistivity with increasing temperature. One can see that the  $\rho(T)$  curves measured for U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub> and UNi<sub>2</sub>Si<sub>2</sub> are qualitatively very similar. The only differences are in the location of the two above mentioned  $\rho(T)$  anomalies, which reflects a slight increase of  $T_2$  and a stronger decrease of  $T_1$  observed with the Pd doping in UNi<sub>2</sub>Si<sub>2</sub>.

The obtained magnetization and resistivity results provide a strong evidence for nearly identical magnetic phase diagrams for  $U(Ni_{0.95}Pd_{0.05})_2Si_2$  and  $UNi_2Si_2$ , *i.e.* the two compounds condense in the same ground state characterized by the UAF magnetic structure although the temperature interval of the stability of the AF-I phase becomes somewhat expanded with Pd doping, especially on the low-temperature side.

# $3.2 U(Ni_{0.91}Pd_{0.09})_2Si_2$

When the Pd content is increased, dramatic changes of magnetism can be recognized. The most striking change is seen in the temperature dependence of resistivity where one observes no anomaly below 50 K, which is accompanying the UAF  $\iff$  AF-I magnetic phase transition in U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub> and UNi<sub>2</sub>Si<sub>2</sub>. On the other hand,



Fig. 4. Comparison of the virgin magnetization curves and hysteresis loops at 4.2 K for compounds with x = 0.05, 0.09 and 0.135.

the resistivity behavior of  $U(Ni_{0.91}Pd_{0.09})_2Si_2$  is reminiscent of  $UPd_2Si_2$  [6]. This would, however, suggest the entire AF-I structure in the ground state exists also for  $U(Ni_{0.91}Pd_{0.09})_2Si_2$ .

On the other hand, the temperature dependences of magnetization (both for the ZFC and FC curves) measured in 0.01 T (see Fig. 2b) exhibit anomalies below 50 K, which are qualitatively similar to those observed for U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub> and UNi<sub>2</sub>Si<sub>2</sub> [2]. The low-temperature FC magnetization value is, however, 5 times lower than that for U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub>. The M(T) dependence in the vicinity of  $T_2$  and  $T_N$  is very similar to that for x = 0.05 and is not shown in Figure 2. The positions of the two magnetization maxima in the ZFC curve reflect a slight increase (decrease) of  $T_2$  ( $T_1$ ) observed with increasing Pd content, which points to stabilization of the AF-I phase. Also the  $T_N$  value is somewhat higher with respect to that of the 5% -Pd compound. Similar conclusions can be drawn from the analysis of resistivity data.

the magnetization curves measured for Also  $U(Ni_{0.91}Pd_{0.09})_2Si_2$  at temperatures up to 40 K are different from corresponding data obtained for the sample with lower Pd concentration. At 4.2 K, two successive magnetization steps around  $0.5(=B_{c1})$  and  $4.5 \text{ T} (=B_{c2})$ are seen in the virgin curve (Fig. 4). The hysteresis loop, however, resembles strongly  $U(Ni_{0.95}Pd_{0.05})_2Si_2$  and  $UNi_2Si_2$  [2]. The remanent magnetic moment equals to  $0.53 \ \mu_{\rm B}/{\rm f.u.}$  Further M(B) cycling repeats the hysteresis loop and the heating up to  $\sim 50$  K is needed to restore the initial state. This magnetization behavior serves an outstanding example of a material, which exhibits the virgin curve being mostly outside the hysteresis loop. In our understanding, the steps on the virgin magnetization curve reflect field induced magnetic phase transitions. The type of hysteresis loop confirms that we deal with an irreversible transition from the AF-I to the UAF phase. In loose words we may say that  $U(Ni_{0.91}Pd_{0.09})_2Si_2$ at 4.2 K represents a rare case of a permanent mag-



Fig. 5. The magnetization isotherms along the *c*-axis of the  $U(Ni_{0.91}Pd_{0.09})_2Si_2$  crystal at selected temperatures.

net (ferromagnet) induced by magnetic field from an antiferromagnet.

At 40 K a two-step metamagnetic transition is observed in the magnetization curve (Fig. 5). In this case the values of critical fields are  $B_{c1} \simeq 0.1$  T and  $B_{c2} \simeq 1.5$  T. The upper transition shows a large hysteresis of approximately 1 T. Except for the small spontaneous magnetization component (~ 0.1  $\mu_{\rm B}/{\rm f.u.}$ ) the 40 K magnetization curves can be attributed to metamagnetic processes in an antiferromagnet. The spontaneous magnetization component vanishes above 40 K and the magnetization curves at higher temperatures resemble the behavior of the compounds with lower Pd content. The metamagnetic transition is, however, shifted to higher (lower) fields with increasing Pd concentration in the interval of 80 < T < 100 K (100 < T < 120 K). This result suggests that the Pd substitution for Ni stabilizes the AF-I phase and destabilizes the LSDW phase with respect to the application of magnetic field.

#### $3.3 U(Ni_{0.865}Pd_{0.135})_2Si_2$

Electrical resistivity data obtained for this composition (Fig. 3) strongly resemble results observed for  $U(Ni_{0.91}Pd_{0.09})_2Si_2$ , only the  $T_2$ -anomaly is shifted to slightly higher temperatures. The temperature dependence of magnetization seen in Figure 2b also shows features analogous to these observed for the compounds with lower Pd content. The low temperature magnetization measured in 0.01 T is reduced in this case down to about 5% of the value observed for  $U(Ni_{0.95}Pd_{0.05})_2Si_2$ .

The M(B) curves (Figs. 4, 6) in many aspects resemble the corresponding magnetization data obtained for the compound with x = 0.09. At 4.2 K, a two-step metamagnetic transition is observed with  $B_{c1} \sim 0.5$  T. The upper magnetization step becomes broader and is shifted to slightly higher magnetic fields (Fig. 4). With increasing temperature, the evolution of magnetization curves



Fig. 6. The magnetization isotherms along the *c*-axis of the  $U(Ni_{0.865}Pd_{0.135})_2Si_2$  crystal at selected temperatures in the AF-I range (a) and in the LSDW range (b).



Fig. 7. Details of low temperature hysteresis loops of  $U(Ni_{0.865}Pd_{0.135})_2Si_2$ .

is very similar to that observed for  $U(Ni_{0.91}Pd_{0.09})_2Si_2$ (Fig. 6). The close similarity of the 4.2 K magnetization curves measured for the two compounds shows that the conclusion about the irreversible AF-I  $\rightarrow$  UAF transition drawn above for  $U(Ni_{0.91}Pd_{0.09})_2Si_2$  is without doubt valid also for  $U(Ni_{0.865}Pd_{0.135})_2Si_2$ . With increasing temperature, the irreversibility becomes less pronounced; a part of volume transfers back to the AF-I state. Figure 7 shows details of hysteresis loops in  $U(Ni_{0.865}Pd_{0.135})_2Si_2$ below 40 K. In a wide field interval, the AF-I and UAF phases coexist which results in the plateau in the M(B)curve. The decreasing height of plateau shows how the AF-I I volume fraction grows on account of UAF with heating.

The low-temperature magnetization values observed for  $U(Ni_{0.91}Pd_{0.09})_2Si_2$  and  $U(Ni_{0.865}Pd_{0.135})_2Si_2$  in 0.01 T represent only small fractions of those observed for

U(Ni<sub>0.95</sub>Pd<sub>0.05</sub>)<sub>2</sub>Si<sub>2</sub>. This result may be interpreted either in terms of emerging new "UAF" phases with very complicated stacking that would yield a spontaneous moment of approximately 0.08  $\mu_{\rm B}/f.u.$  and 0.02  $\mu_{\rm B}/f.u.$  This could be achieved only by a very complex stacking of U moments leading to a very large magnetic elementary cell, which would imply observation of a strong increase of the electrical resistivity at  $T_1$  with decreasing temperature, which is apparently not the case. In this model also the state at the plateau in the M(B) curves in Figure 7 would be considered as a homogeneous state with complex stacking of U moments instead of a mixture of AF-I and UAF phases.

Another scenario may be formulated in terms of the above discussed coexistence of the AF-I and UAF phases at temperatures below  $T_1$  with fractional volumes depending on the composition of the Ni/Pd sublattice and on external parameters (magnetic field, temperature or pressure). The dramatic reduction of the low temperature FC (in 0.01 T) magnetization with Pd doping should then reflect the rapidly decaying fractional volume of the UAF phase to become less than 4% of the total volume in  $U(Ni_{0.865}Pd_{0.135})_2Si_2$ . This scenario could explain also the absence of resistivity anomaly at  $T_1$  both for the  $U(Ni_{0.91}Pd_{0.09})_2Si_2$  and  $U(Ni_{0.865}Pd_{0.135})_2Si_2$  compounds. The majority volume with the AF-I characterized by the lower resistivity in these two compounds by passes the more resistive UAF clusters, which are emerging below  $T_1$ . Nevertheless, the possible appearance of a new magnetic phase with different propagation vector cannot be entirely ruled out. Therefore, the microscopic experiments, especially the neutron diffraction, are strongly desired in order to confirm the ground state of the x = 0.09and 0.135 samples.

We measured also the temperature dependence of susceptibility at temperatures between  $T_{\rm N}$  and room temperature in a field of 4 T applied along the c-axis. In case of all our samples and also  $UNi_2Si_2$  and  $UPd_2Si_2$ the susceptibility follows well the Curie-Weiss law. The values of  $\Theta_p$  vary from -4 K in UNi<sub>2</sub>Si<sub>2</sub> [10] to 2 K for  $U(Ni_{0.865}Pd_{0.135})_2Si_2$ . A linear extrapolation of this trend to  $UPd_2Si_2$  yields 36 K, which compares very well with 40 K [6]. The values of the effective moment  $\mu_{\text{eff}}$ for  $U(Ni_{1-x}Pd_x)_2Si_2$  appear in the narrow interval between 3.4 and 3.5  $\mu_{\rm B}/{\rm U}$  atom. (3.34  $\mu_{\rm B}$  for UNi<sub>2</sub>Si<sub>2</sub> and  $3.5~\mu_{\rm B}$  for UPd\_2Si\_2). The ordered-moment value does not show any considerable dependence on the Pd concentration ( $\mu_{\rm ord} = 0.55 \pm 0.02 \ \mu_{\rm B}$ ). This reflects the stability of the U magnetic moment with respect to the varying chemical neighborhood of U atoms whereas the exchange interactions between the U moments is somewhat sensitive to the changes in the Ni/Pd sublattice.

# 4 Conclusions

In conclusion, we have prepared single crystals of  $U(Ni_{1-x}Pd_x)_2Si_2$  compounds with x = 0.05, 0.09 and 0.135. Magnetization isotherms at different temperatures and temperature dependences of magnetic susceptibility

and electrical resistivity have been measured. The magnetic properties of x = 0.05 sample is similar to that of pure UNi<sub>2</sub>Si<sub>2</sub> and it seems to have entirely the UAF ground state. On the other hand, in the ground state of the samples with x = 0.09 and 0.135, very small magnetic moment has been observed in the ZFC sample at low temperatures. Moreover the increase of electrical resistivity accompanied by a magnetic order-order transition from AF-1 to UAF has not been detected for  $x \ge 0.09$ . Application of a sufficiently high magnetic field along the *c*-axis causes an irreversible transition to the UAF phase at lowest temperatures. The UAF phase persists at low temperatures after removing external magnetic field. These phenomena can be attributed to the ground state of x = 0.09and 0.135 samples being characterized by the coexistence an AF-I and UAF with a small volume fraction of the UAF phase. In the mean time between submission of this paper and final acceptance we performed a neutron diffraction study of the two U(Ni<sub>1-x</sub>Pd<sub>x</sub>)<sub>2</sub>Si<sub>2</sub> single crystals with x = 0.09 and 0.135 [11]. We have confirmed validity of the proposed scenario including the fact that the Néel temperature  $(T_N)$  of both these compounds is almost the same as in  $\text{UNi}_2\text{Si}_2$ , *i.e.*  $\simeq 125$  K, because clear magnetic neutron diffraction reflections relevant to AF ordering with q =(0,0,0.76) have been observed below this temperature.

This work is a part of the research program MSM113200002 that is financed by the Ministry of Education of the Czech

Republic. It has been partially supported (F.H.) also by the Grant Agency of the Czech Republic (grants # 202/01/D045).

### References

- Pearson's Handbook of Crystallographic Data for Intermetallic Phases, 2nd edn., edited by P. Villars, L.D. Calvert (American Society for Metals, Materials Park, OH, 1991)
- V. Sechovský, L. Havela, in *Handbook of Magnetic Materials*, edited by K.H.J. Buschow, Vol. 11 (Elsevier Science B.V., Amsterdam, 1998), p. 1, and references therein
- B. Shemirani, H. Lin, M.F. Collins, C.V. Stager, J.D. Garrett, W.J.L. Buyers, Phys. Rev. B 47, 8672 (1993)
- H. Lin, L. Rebelsky, M.F. Collins, J.D. Garrett, W.J.L. Buyers, Phys. Rev. B 43, 3232 (1991)
- 5. P. Svoboda, P. Javorsky, V. Sechovský, A.A. Menovsky, Physica B **276**, 686 (2000)
- T. Honma, H. Amitsuka, T. Sakakibara, K. Sugiyama, M. Date, Physica B 186-188, 684 (1993)
- Y.B. Ning, J.D. Garrett, W.R. Datars, J. Phys. Cond. Matt. 4, 9995 (1992)
- M.M. Barati, W.R. Datars, T.R. Chien, C.V. Stager, J.D. Garrett, Phys. Rev. B 48, 16926 (1993)
- J. Vejpravová, P. Svoboda, V. Sechovský, C. Ritter, Appl. Phys. A (in press), (2002)
- T. Takeuchi, K. Watanabe, T. Taniguchi, T. Kuwai, A. Yamagishi, Y. Miyako, Physica B 201, 243 (1994)
- P. Svoboda, F. Honda, A.V. Andreev, J. Vejpravová, K. Prokeš, V. Sechovský, Physica B (in press), (2003)